

A solution to fractal Hilbert's 16th problem for slow-fast Liénard equations

Renato Huzak (Hasselt University, Belgium)

joint work with P. De Maesschalck, A. Janssens and G. Radunovic (FRABDYN HRZZ PZS 3055)

10.5.2023



Traditional Hilbert's 16th problem (difficult) —> Fractal Hilbert's 16th problem

Traditional Hilbert's 16th problem (difficult) \longrightarrow Fractal Hilbert's 16th problem

Slow-fast Liénard equations

$$\dot{x} = y - F_{n+1}(x), \quad \dot{y} = -\epsilon G_m(x)$$

Traditional Hilbert's 16th problem

$$\dot{x} = P_n(x, y), \quad \dot{y} = Q_n(x, y)$$

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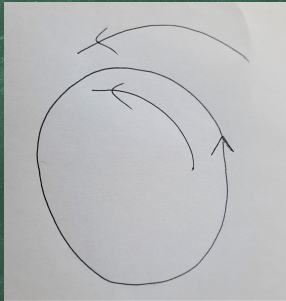
Is there a bound $H(n)$ on the number of limit cycles?

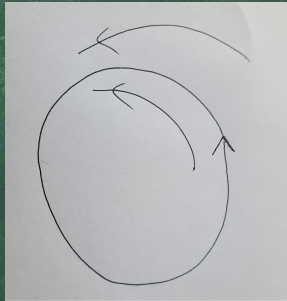
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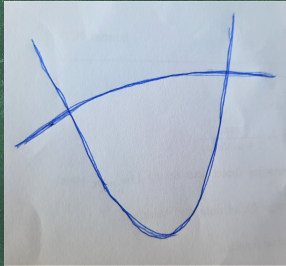
Is there a bound $H(n)$ on the number of limit cycles?

Is $H(n) \leq n^a$? (Smale)

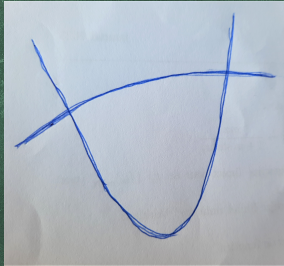




$H(2) < \infty$? Dumortier-Roussarie-Rousseau program, 1994



solve it locally in the (x, y) -phase space and parameter space



solve it locally in the (x, y) -phase space and parameter space

Our goal in slow-fast setting: the number of limit cycles \longrightarrow
cardinality of the fractal spectrum

Slow-fast Liénard equations

$$\begin{cases} \dot{x} = y - \sum_{k=0}^{n+1} B_k x^k & A_m \neq 0 \\ \dot{y} = -\epsilon \sum_{k=0}^m A_k x^k & B_{n+1} \neq 0 \end{cases}$$

$\epsilon \geq 0$ is the singular perturbation parameter kept small

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$m = 1$: classical Liénard equations of degree $n + 1$

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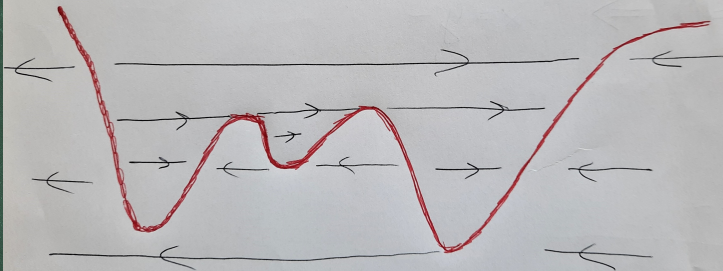
$\epsilon \geq 0$ is the singular perturbation parameter kept small

$m = 1$: classical Liénard equations of degree $n + 1$

$m > 1$: generalized Liénard equations

$n = 1$: Liénard equations with linear damping

$$\varepsilon = 0 \Rightarrow \begin{cases} \dot{x} = y - \sum_{k=0}^{n+1} B_k x^k \\ \dot{y} = 0 \end{cases} \quad \leftarrow \text{Fast Subsystem}$$



$$\left\{ \begin{array}{l} \dot{x} = y - \sum_{k=0}^{n+1} B_k x^k \quad A_m \neq 0 \\ \dot{y} = -\varepsilon \sum_{k=0}^m A_k x^k \quad B_{n+1} \neq 0 \end{array} \right.$$

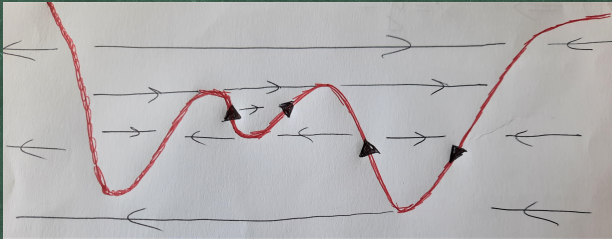
$$\begin{cases} \dot{x} = y - \sum_{k=0}^{n+1} B_k x^k & A_m \neq 0 \\ \dot{y} = -\varepsilon \sum_{k=0}^m A_k x^k & B_{n+1} \neq 0 \end{cases}$$

$$\begin{cases} \varepsilon x' = y - \sum_{k=0}^{n+1} B_k x^k \\ y' = -\sum_{k=0}^m A_k x^k \end{cases} \xrightarrow{\varepsilon \rightarrow 0} \begin{cases} 0 = y - \sum_{k=0}^{n+1} B_k x^k \\ y' = -\sum_{k=0}^m A_k x^k \end{cases}$$

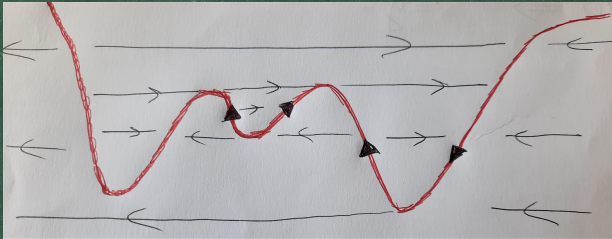
Slow Subsystem

$$\Rightarrow x' = -\frac{\sum_{k=0}^m A_k x^k}{\left(\sum_{k=0}^{n+1} B_k x^k\right)'}$$

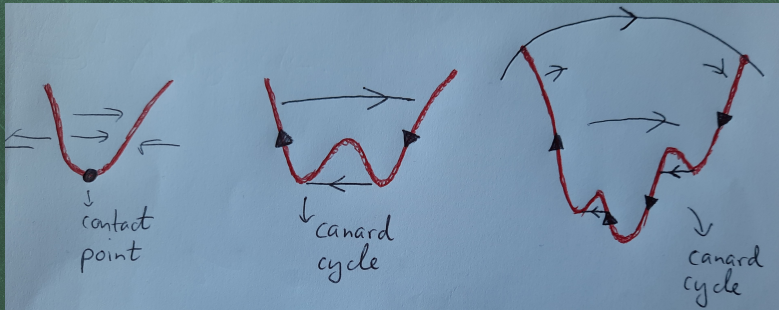
slow+fast



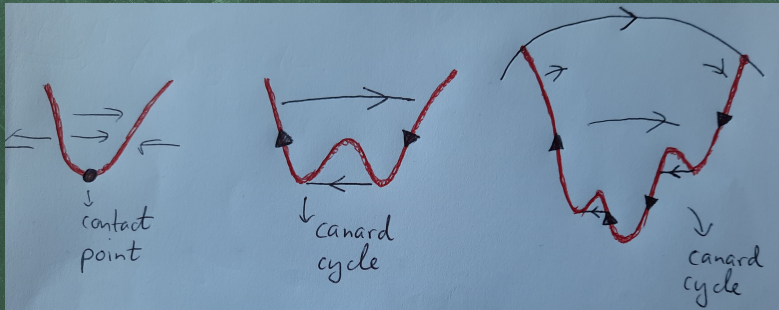
slow+fast



Limit periodic sets

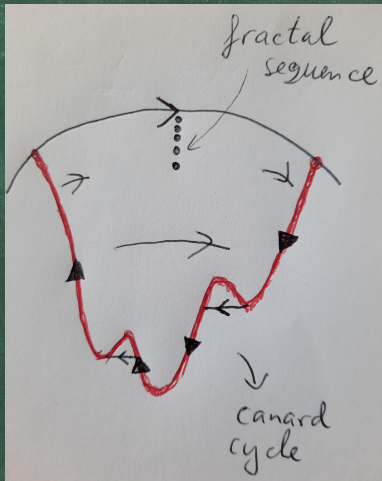


Limit periodic sets



Canard cycles... [De Maesschalck, Dumortier, Roussarie, 2021]

sequences near infinity



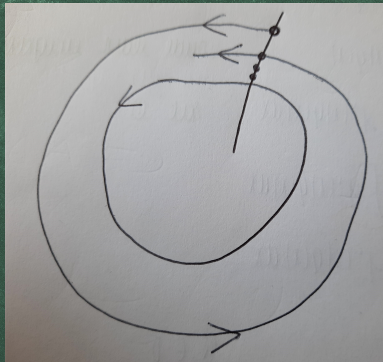
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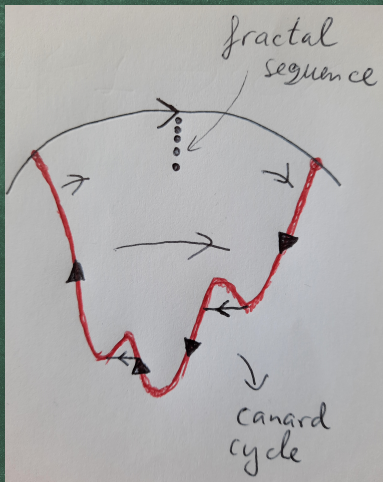
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regular case [Zubrinic,Zupanovic,2005,2008]



to generate a sequence we use entry-exit (or slow-relation function)



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Entry-Exit

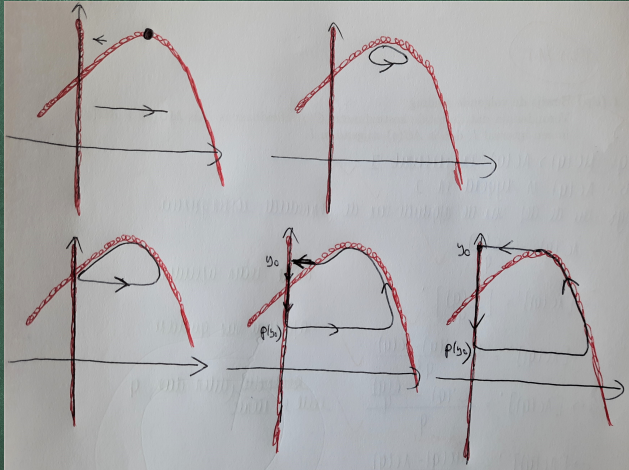
$$\begin{cases} \dot{x} = \bar{f}(x, y, \varepsilon) x^n \\ \dot{y} = \varepsilon \bar{g}(x, y, \varepsilon) \end{cases} \quad n \geq 1$$

- $\bar{g}(0, y, 0) < 0$
- $\bar{f}(0, y, 0) > 0$ for $y < 0$
- $\bar{f}(0, y, 0) < 0$ for $y > 0$

ENTRY-EXIT:

$$P(y_0) = \int_{y_0}^{\infty} \frac{\bar{f}(0, y, 0)}{\bar{g}(0, y, 0)} dy = 0$$

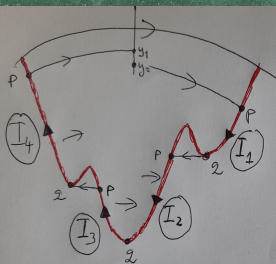
Predator-prey systems



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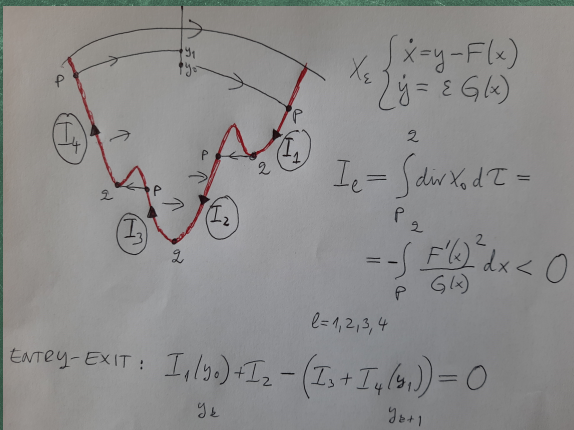
$$X_\varepsilon \begin{cases} \dot{x} = y - F(x) \\ \dot{y} = \varepsilon G(x) \end{cases}$$

$$I_\varepsilon = \int_{P_2}^2 \text{div} X_0 d\mathcal{T} =$$

$$= - \int_P^2 \frac{F'(x)^2}{G(x)} dx < 0$$

$$l=1,2,3,4$$

$$\text{ENTRY-EXIT: } \underbrace{I_1(y_0)}_{y_\varepsilon} + I_2 - \left(I_3 + \underbrace{I_4(y_1)}_{y_{\varepsilon+1}} \right) = 0$$



$$I(y) \stackrel{\text{def}}{=} I_1(y) + I_2 - (I_3 + I_4(y)) \rightarrow \text{slow div. integral}$$

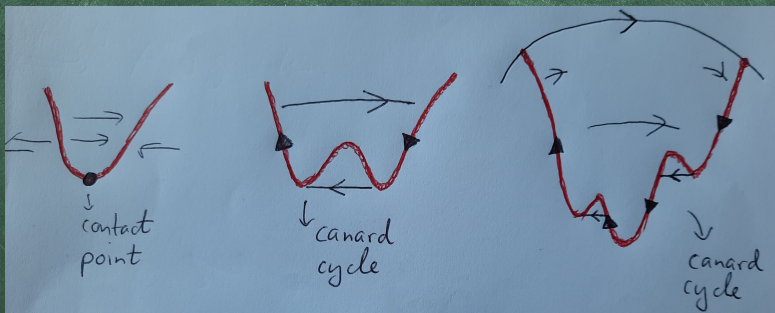
$$I(y) \rightarrow \begin{matrix} I_* \\ \neq 0 \\ 0 \end{matrix}, 0, \pm\infty \text{ as } y \rightarrow +\infty$$

$$\frac{F'(x)^2}{G(x)} = \alpha x^{2n-m} (1 + o(1)), \quad x \rightarrow \pm\infty$$

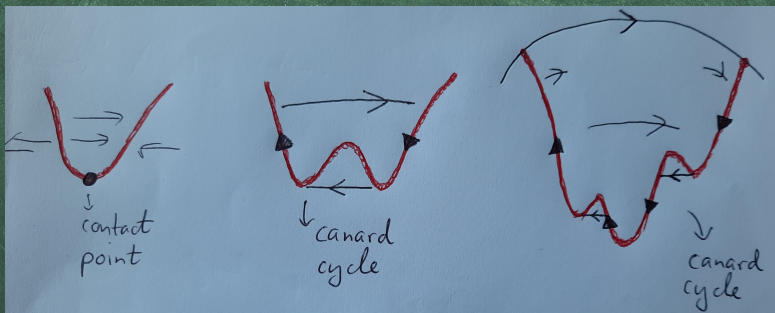
$$m < 2n+1$$

$$m = 2n+1$$

$$m > 2n+1$$



[Huzak,2017] [Huzak,Vlah,2019] [Crnkovic,Huzak,Vlah,2021]
 [Huzak,Vlah,Zubrinic,Zupanovic,2023] [De
 Maesschalck,Huzak,Janssens,Radunovic,2023]



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 Maesschalck,Huzak,Janssens,Radunovic,2023]
 [Dmitrovic, Huzak, Vlah, Zupanovic,2021],
 [Huzak,Mardesic,Resman,Zupanovic,2023],
 [Crnkovic,Huzak,Resman,2023]

Box dimension (see Falconer, Lapidus, Tricot, ...)

- Let $\delta > 0$ and $\delta \approx 0$
- $U(\delta) \stackrel{\text{def}}{=} \text{the } \delta\text{-neighborhood of a bounded } U \subseteq \mathbb{R}$
- $|U(\delta)| \stackrel{\text{def}}{=} \text{the Leb. measure of } U(\delta)$

\Rightarrow the lower box dimension:

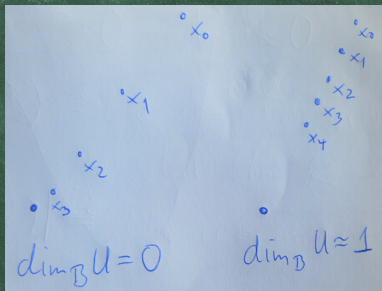
$$\underline{\dim}_{\mathbb{B}} U = \liminf_{\delta \rightarrow 0} \left(1 - \frac{\ln |U(\delta)|}{\ln \delta} \right)$$

the upper box dimension:

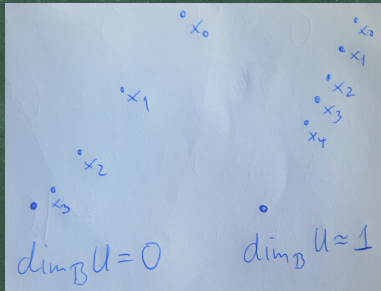
$$\overline{\dim}_{\mathbb{B}} U = \limsup_{\delta \rightarrow 0} \left(1 - \frac{\ln |U(\delta)|}{\ln \delta} \right)$$

$\dim_{\mathbb{B}} U$

Box dimension measures the density of U —the bigger the box dimension, the higher the density



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See [Box dimension of trajectories of 1-dim discrete dynamical systems, Elezovic, Zupanovic, Zubrinic, 2007]

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$m \leq 2n + 1$: Poincaré-Lyapunov disc of degree $(1, n + 1)$

$$x = \bar{x}/r, y = 1/r^{n+1}$$

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$$x = \pm 1/r, y = \bar{y}/r^{\frac{m+1}{2}}$$

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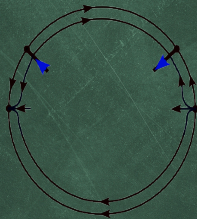
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—a rescaling in x, y, t —

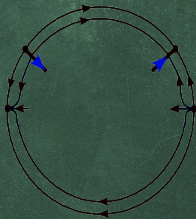
$$\begin{cases} \dot{x} = y - \left(x^{n+1} + \sum_{k=0}^n b_k x^k \right) \\ \dot{y} = -\varepsilon \left(A x^m + \sum_{k=0}^{m-1} a_k x^k \right) \end{cases}$$

$A = 1$ if m is even and $m \neq 2n+1$
 $A = \text{sign}(A_m)$ if m is odd and $m \neq 2n+1$
 $A \neq 0$ if $m = 2n+1$

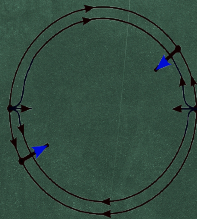
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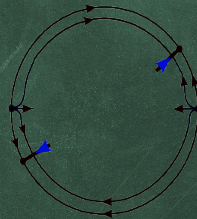
$A > 0, n \text{ odd}$



$A < 0, n \text{ odd}$

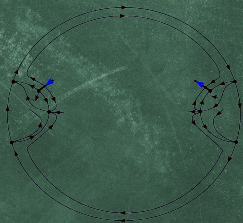


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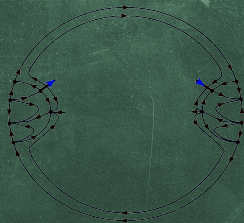


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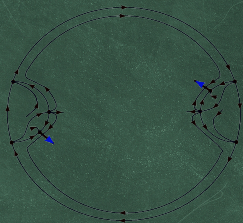
$m > 2n + 1$, m odd



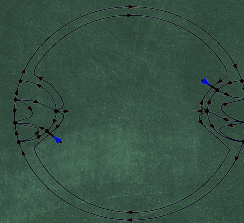
$A = 1$, n odd



$A = -1$, n odd

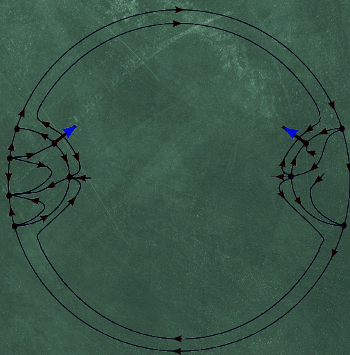


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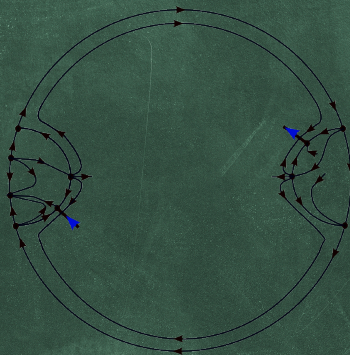


$A = -1$, n even

$m > 2n + 1$, m even



$A = 1$, n odd



$A = 1$, n even

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$A = \text{sign}(A_m)$ if m is odd and $m \neq 2n+1$

$A \neq 0$ if $m = 2n+1$

we break the symmetry: $b_{2j+1} = a_{2k} = 0$

Theorem ($m < 2n + 1$)

$$\dim_{\mathbb{B}} U = \frac{n-2j}{n+1-2j}$$

$j = 0, \dots, \frac{n-1}{2}$

$I(y) \rightarrow \pm\infty$
or
 $I(y) \rightarrow 0$

$$\dim_{\mathbb{B}} U = \frac{m-2k}{m+1-2k}$$

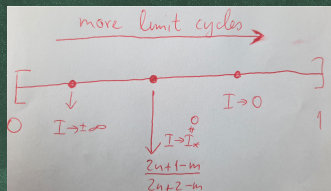
$k = 0, \dots, \frac{m-1}{2}$

$I(y) \rightarrow \pm\infty$
 $I(y) \rightarrow 0$

$$\dim_{\mathbb{B}} U = \frac{2n+1-m}{2n+2-m}$$

$I(y) \rightarrow I_* \in \mathbb{R}$

0



Theorem ($m = 2n + 1$)

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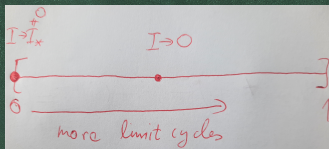
$$I(y) \rightarrow 0$$

$$\dim_{\mathbb{B}} U = \frac{2n+1-2k}{2n+2-2k}$$

$$I(y) \rightarrow 0$$

$$\dim_{\mathbb{B}} U = 0$$

$$I(y) \rightarrow \underset{\substack{+ \\ 0}}{I_*} \in \mathbb{R}$$

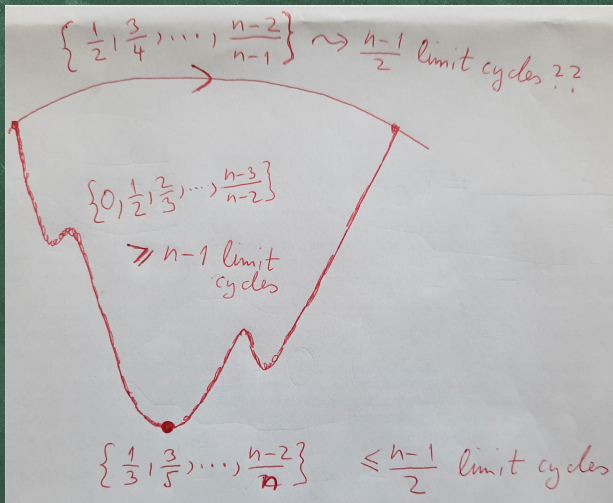


Theorem ($m > 2n + 1$)

$I(y) \rightarrow 0$

$$\dim_{\mathbb{B}} U = \frac{(n-2j)(m+1)}{(n-2j)(m+1) + 2(n+1)}$$
$$\dim_{\mathbb{B}} U = \frac{(m-2k)(m+1)}{(m-2k)(m+1) + 2(n+1)}$$

Classical Liénard equations ($m < 2n + 1$)



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